Calibration II

Dharam V. Lal

with due thanks to several friends / collaborators at

IDIA-UCT (SA), NCRA-TIFR (India) and NRAO (USA)

CALIBRATION AND IMAGING

Standard calibration and imaging (DI instrumental effects) w/ DD instrumental + propagation effects correction for *w*-term and for PB image plane correction Fourier plane correction pointing self-calibration Mosaicing w/ advanced image parameterisation multi-scale CLEAN (deconvolution) multi-frequency synthesis (imaging) full polarisation (Stokes) calibration and imaging

TELESCOPE SENSITIVITY

Noise limit for imaging with interferometric radio telescopes

$$
\sigma = \frac{T_{\rm sys}}{A_{\rm eff} \times \sqrt{(\Delta \nu \times \Delta t)}}
$$

Sensitivity improvements achieved by

- wide band receivers,
- long integration times
- more antennas
- long baselines

$$
\sigma_{\text{confusion}} \propto (\nu^{-2.7}/B_{\text{max}}^2)
$$

 $B_{\text{max}} \sim 100 \text{ km } \omega$ 200 MHz, the confusion noise is \sim 1 μ Jy beam⁻¹.

IMAGING CHALLENGES AT LOW FREQ.

- Wide-field imaging
	- account for direction dependent (DD) effects
	- PB: time, frequency and polarisation dependence

w-term

Wide-band imaging

… plus frequency dependence of the sky brightness Data volume $\propto N_{\text{ant}}^2 \times N_{\text{channel}} \times t$ Sky brightness \Longrightarrow multi-scale deconvolution Ionospheric effects \Longrightarrow need for DD solvers

$$
\overrightarrow{V}_{ij}^{Obs}(\nu, t) = G_{ij} W_{ij}(\nu, t) \int P_{ij} M_{ij}(s, \nu, t) \overrightarrow{I}(s, \nu) e^{i\overrightarrow{b}_{ij} \cdot \overrightarrow{s}} d\overrightarrow{s}
$$

$$
V(u, v, w) = \int \frac{I(l, m, n)}{n} e^{-2\pi i (ul + v m + w(n-1))} dl dm
$$

$$
\Gamma(\vec{b}) = \int \left\langle E(\vec{s}, t) \cdot E^*(\vec{s}, t - \vec{b} \cdot \vec{s}/c) \right\rangle e^{-2\pi i \vec{b} \cdot \vec{s}/\lambda} d\Omega
$$

\n
$$
V(u, v, w) = \int \frac{I(l, m, n)}{n} e^{-2\pi i (ul + v m + w(n-1))} dl dm
$$

\n(for $w \approx 0, n \approx 1$)
\n
$$
V(u, v) = \int I(l, m) e^{-2\pi i (ul + v m)} dl dm
$$

\n(this is you result?

(this is van-Cittert Zernike theorem)

$$
\Gamma(\vec{b}) = \int \left\langle E(\vec{s}, t) \cdot E^*(\vec{s}, t - \vec{b} \cdot \vec{s}/c) \right\rangle e^{-2\pi i \vec{b} \cdot \vec{s}/\lambda} d\Omega
$$

$$
V(u, v, w) = \int \frac{I(l, m, n)}{n} e^{-2\pi i (ul + v m + w(n-1))} dl dm
$$

Polarised radiation:

$$
\overrightarrow{E}_{i} = [E^{r} E^{l}]_{i}^{T}
$$

\n(four cross-correlation products, $\langle \overrightarrow{E}_{i} \otimes \overrightarrow{E_{j}^{*}} \rangle$ per baseline)
\n
$$
\overrightarrow{V}_{ij} = [V^{rr} V^{rl} V^{lr} V^{ll}]_{ij}^{T}
$$
\n
$$
\overrightarrow{I} = [I^{rr} I^{rl} I^{lr} I^{ll}]^{T}
$$
\n
$$
\overrightarrow{R}
$$
\n
$$
\downarrow
$$
\n
$$
\overrightarrow{C
$$
\n
$$
\downarrow
$$

 $\overrightarrow{E_i} = [E^r \ E^l]$ *T i*

(suffers from propagate effects and receiver electronics) (Jones matrices describe this modulation for the electric field incident at each orthogonal pair of feeds $\overrightarrow{E}_i = [E^r \ E^l]_i^T$ $D I: J_i^{vis} = [G D C]$ $(a 2 \times 2$ matrix product) complex gains, *G*, polar'n leakage, *D* and feed config'n, *C. T i* DD: *Jsky* $(a 2 \times 2$ matrix product) AIPs, *E*, PA effects, *P* and tropospheric / ionospheric effects, and Faraday R'n, *F*. *i* = [*EPF*]

 $\overrightarrow{E_i} = [E^r \ E^l]$ *T i*

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(Jones matrices describe this modulation for the electric field incident at each orthogonal pair of feeds $\overrightarrow{E}_i = [E^r \ E^l]_i^T$ *T i*

DI:
$$
J_i^{vis} = [GDC]
$$

\nDD: $J_i^{sky} = [EPF]$

\n $K_{ij}^{\{vis, sky\}} = [J_i \otimes J_j^{\dagger}]^{\{vis, sky\}}$

 $I \cap I \cap I$

(effect on each baseline *ij* is described by the outer-product of these antenna-based Jones matrices, a 4×4 matrix!)

$$
\overrightarrow{V}_{ij}^{obs} = [K_{ij}^{vis}] \int [K_{ij}^{sky}] \overrightarrow{I}^{sky}(\overrightarrow{s}) e^{-2\pi i \overrightarrow{b} \cdot \overrightarrow{\sigma} / \lambda} d\Omega
$$

$$
\overrightarrow{V}_{ij}^{obs} = [K_{ij}^{vis}] \left[[K_{ij}^{sky}] \overrightarrow{I}^{sky}(\overrightarrow{s}) e^{-2\pi i \overrightarrow{b} \cdot \overrightarrow{\sigma} / \lambda} d\Omega \right]
$$

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Standard calibration and imaging (DI instrumental effects) w/ DD instrumental + propagation effects correction for *w***-term** and **for PB image plane** correction Fourier plane correction pointing self-calibration Mosaicing w/ advanced image parameterisation **multi-scale CLEAN** (deconvolution) **multi-frequency synthesis** (imaging) **full polarisation** (Stokes) calibration and imaging

W-TERM

 $V(u, v, w) =$ $e^{iw\sqrt{1-l^2-m^2}}$

divide the FoV into a no. of FACETS

n

Credits: S. Bhatnagar, synthesis imaging NRAO workshop

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W-TERM

$$
V(u, v, w) = \int \frac{I(l, m, n)}{n} e^{-2\pi i (ul + v m + w(n-1))} d
$$

$$
K_{ij}^{Sky} = e^{w_{ij} (\sqrt{1 - l^2 - m^2} - 1)}
$$

An order-of magnitude faster than FACETing, and

for the same amount of computing time provides higher DR images.

Credits: S. Bhatnagar, synthesis imaging NRAO workshop

A-projection

A-projection

Visibility depends on time and frequency!

CALIBRATION AND IMAGING

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multi-frequency synthesis

$$
I_{\nu}^{sky} = I_{\nu 0}^{sky} \left(\frac{\nu}{\nu 0}\right)^{I_{\alpha}^{sky} + I_{\beta}^{sky} log(\frac{\nu}{\nu 0})}
$$

A-projection

 $\overrightarrow{V}_{ij}^{obs}(\nu) = \int J_{ij}^{S}$ *ij* (*s*, *ν*, *t*) *I*(*s*, *ν*) *ei*(*ul*+*vm*+*w*((1−*^l* 2 −*m*²)−1)) *ds*

A-projection

$$
\overrightarrow{V}_{ij}^{obs}(\nu) = \int J_{ij}^{S}(s, \nu, t) I(s, \nu) e^{i(ul+vm+w(\sqrt{(1-l^2-m^2)-1}))} ds
$$

multi-frequency synthesis

$$
I_{\nu}^{sky} = I_{\nu 0}^{sky} \left(\frac{\nu}{\nu 0}\right)^{I_{\alpha}^{sky} + I_{\beta}^{sky} log(\frac{\nu}{\nu 0})}
$$

$$
I_0 = I_{\nu_0}
$$

$$
I_1 = I_{\alpha} \times I_{\nu_0}
$$

$$
I_2 = (I_{\alpha}(I_{\alpha} - 1)/2 + I_{\beta}) \times I_{\nu_0}
$$

- 31 **140 MHz Dharam V. LAL (NCRA-TIFR)**

RAS2024: 26 Nov 2024

CORRECTION FOR **PB**

FT

(standard imaging)

FT + MT-MFS

+ A-projection

FT + MT-MFS + WB A-projection

Bhatnagar (NRAO,

 $\dot{\mathbf{v}}$

Credits:

PEELING: DD CALIBRATION

antenna based gains are determined in the direction of each compact source.

subtract these gains (contribution of compact sources from the observed data using a DFT) and the residual visibilities are imaged again.

drawbacks of peeling…

PEELING: DD CALIBRATION

antenna based gains are determined in the direction of each compact source.

subtract these gains (contribution of compact sources from the observed data using a DFT) and the residual visibilities are imaged again. drawbacks of peeling…

Credits: H. Intema (Leiden Obs.)

MORE DD ISSUES: CORRECTION FOR **PB**

A-projection

$$
\overrightarrow{V}_{ij}^{obs} = [K_{ij}^{vis}] \int [K_{ij}^{sky}] \overrightarrow{I}^{sky}(\overrightarrow{s}) e^{-2\pi i \overrightarrow{b} \cdot \overrightarrow{\sigma} / \lambda} d\Omega
$$

Remember, 200 / 400 MHz bandwidths and
Sky is (not) variable, and
Antenna power pattern is (not) $\alpha^{\frac{2}{8}}$ Assumption: sky is (not) variable, and Antenna power pattern is (not) $c^{\frac{3}{2}}$

RASA RASA RASA RASA STOKES PARAMETERS

I – total intensity and sum of any two orthogonal polarisations

Q & *U* – completely specify linear polarisation *V* – completely specifies circular polarisation $RR = \mathcal{A}(RR)e^{i\psi RR} = I + V$ $LL = \mathcal{A}(LL)e^{i\psi LL} = I - V$ $RL = \mathcal{A}(RL)e^{i\psi RL} = Q + iU$ $LR = \mathcal{A}(LR)e^{i\psi LR} = Q - iU$

RASA RASA RASA RASA STOKES PARAMETERS

I – total intensity and sum of any two orthogonal polarisations

Q & *U* – completely specify linear polarisation *V* – completely specifies circular polarisation Stokes parameters (as percentages of *I)* $I =$ (*RR* + *LL*) 2 *V RR* − *LL I* = *RR* + *LL Q I* = $Re(RL + LR)$ *RR* + *LL U* Im(*RL* − *LR*) *I* = *RR* + *LL*

RASA RASA RASA RASA STOKES PARAMETERS

I – total intensity and sum of any two orthogonal polarisations

Q & *U* – completely specify linear polarisation

V – completely specifies circular polarisation

Stokes parameters (as percentages of *I)*

Is it really that simple?

No, there are leakages…

The total intensity can leak into the polarised components (*I* into {*Q,U,V*}).

Hans Mueller

MUELLER MATRIX

The leakage of each polarisation into the other can be measured and quantified in a 4×4 matrix (Mueller 1943).

POLARISATION CALIBRATION

Flux density scale

- *I Q* leakage ⇔
- *I* ⇔ *U* leakage
- *I* ⇔ *V* leakage

Alignment => PA calibration

Ellipticity, $Q \Leftrightarrow V$

 RL phase, $U \Leftrightarrow V$

Constrained using calibrator with known Stokes parameters

Need calibrator with known PA

Stokes $V \sim 0$ for most calibrators so no need to worry too much unless you require very high precision

PUTTING THIS ALL TOGETHER

In the end what we are trying to do is relate products from our correlator to the intrinsic polarised radiation from the source.

So we need to correct the raw correlator outputs for

imperfections in the receiver (leakages).

The orientation of the receiver with respect to the telescope structure.

a.k.a. the changing parallactic angle.

Any measured propagation related polarisation effects (e.g. Faraday rotation).

For point sources, all of the previous is fine.

What if the source you are looking at is extended compared to the telescope beam?

There are instrumental beam effects that can confuse the measurement of extended polarised signals. They are…

Squint

Squash

What if the source y compared to the tele measurement of extended Squint

Squash

$$
\overrightarrow{V}_{ij}^{Obs}(\nu, t) = W_{ij}(\nu, t) \int M_{ij}(s, \nu, t) \overrightarrow{I}(s, \nu) e^{i\overrightarrow{b_{ij}}, \overrightarrow{s}} d\overrightarrow{s}
$$
\n
$$
M_{ij}(\overrightarrow{s}, \nu, t) = E_i(\overrightarrow{s}, \nu, t) \otimes E_j^*(\overrightarrow{s}, \nu, t)
$$
\n
$$
\overrightarrow{V}_{ij}^{Obs}(\nu, t) = W_{ij}(\nu, t) \mathcal{F} \left[\left(E_i(\overrightarrow{s}, \nu, t) \otimes E_j^*(\overrightarrow{s}, \nu, t) \right) \cdot \overrightarrow{I}(\overrightarrow{s}, \nu) \right]
$$
\n
$$
= W_{ij}(\nu, t) \left[A_{ij} \star \overrightarrow{V}_{ij} \right]
$$
\nwhere, $A_{ij} = A_i \otimes A_j^*$ and \overrightarrow{P} and \overrightarrow{P} and \overrightarrow{P} and \overrightarrow{P} are independent.

Dharam V. LAL (NCRA-TIFR)

DD & DI EFFECTS

$$
\overrightarrow{V}_{ij}^{Obs}(\nu, t) = W_{ij}(\nu, t) \int M_{ij}(s, \nu, t) \overrightarrow{I}(s, \nu) e^{i\overrightarrow{b}_{ij} \cdot \overrightarrow{s}} d\overrightarrow{s}
$$

$$
M_{ij}(\overrightarrow{s}, \nu, t) = E_i(\overrightarrow{s}, \nu, t) \otimes E_j^*(\overrightarrow{s}, \nu, t)
$$

Jones matrix formulation:

A-TO-Z SOLVER

Use Zernike polynomials to directly model the complex aperture

it is a natural domain to model optical aberrations that cause PB weirdness

(Telescope agnostic - does not require ray traced model for different antennas/telescopes, only Holography)

Aperture size is fixed, independent of number of measured sidelobes.

Sekhar (UCT-IDIA) ⇔⇔Jagannathan (NRAO) Sekhar (UCT-IDIA) **Credits: S. Sekhar (UCT-IDIA)** $\dot{\mathbf{v}}$

 $\dot{\mathbf{v}}$

APERTURE ILLUMINATION PATTERN

DHARAM V. LAL (NCRA)

Credits: P. Jagannathan (NRAO)

 \mathbf{a}

Credits: A.Pal (NCRA-TIFR), S. Sekhar (UCT-IDIA)A.Pal (NCRA-TIFR) Credits:

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 $\bf\hat{\psi}$

ANTENNA HOLOGRAPHY

DD effects: getting the beams

Ref. ant. will track the bright point source $(P(0,0) = 1)$

Target ant. will scan the target, but with an offset, response ∝ *P*(*l*, *m*)

cross-corr. \propto beam response of the target antenna, $P(l, m)$

UGMRT DATA

Holography data

UGMRT DATA

Holography data scans/data as a function of time

UNDERSTANDING DATA

Holography data

scans/data as a function of time

construct (Stokes) image/beam

Understanding the polarisation properties of the GMRT dish is fundamental …

need for

"accurate aperture model" **CASA** implementation

DD EFFECTS: GETTING THE BEAMS

Beams:

Once the beams are available, turn-ON AW-projection in CASA to take care of the beam rotation.

Thank you all for your attention!