Calibration II

Dharam V. Lal

with due thanks to several friends / collaborators at

UCT, IDIA (SA), NCRA-TIFR (India) and NRAO (USA)

TELESCOPE SENSITIVITY

Noise limit for imaging with interferometric radio telescopes

 $\sigma = \frac{T_{\rm sys}}{A_{\rm eff} \times \sqrt{(\Delta \nu \times \Delta t)}}$ Sensitivity improvements achieved by wide band receivers, long integration times more antennas $\sigma_{\rm confusion} \propto (\nu^{-2.7}/B_{\rm max}^2)$ long baselines

 $B_{max} \sim 100 \text{ km} @200 \text{ MHz}$, the confusion noise is ~1 μ Jy beam⁻¹.

IMAGING CHALLENGES AT LOW FREQ.

- Wide-field imaging
 - account for direction dependent (DD) effects
 - PB: time, frequency and polarisation dependence

w-term

Wide-band imaging

... plus frequency dependence of the sky brightness Data volume $\propto N_{ant}^2 \times N_{channel} \times t$ Sky brightness \Longrightarrow multi-scale deconvolution Ionospheric effects \Longrightarrow need for DD solvers

IMPLICATIONS FOR IMAGING

Long baselines $B_{\text{max}} > 2 \text{ km} \implies \text{DR} > 10^4$ Wide-field effects:

w-term, PB effects and ionosphere effects Larger data volume

Wide-field, wide-band, high resolution, high dynamic range imaging using large data sizes

a natural consequence of low frequency and high sensitivity imaging.

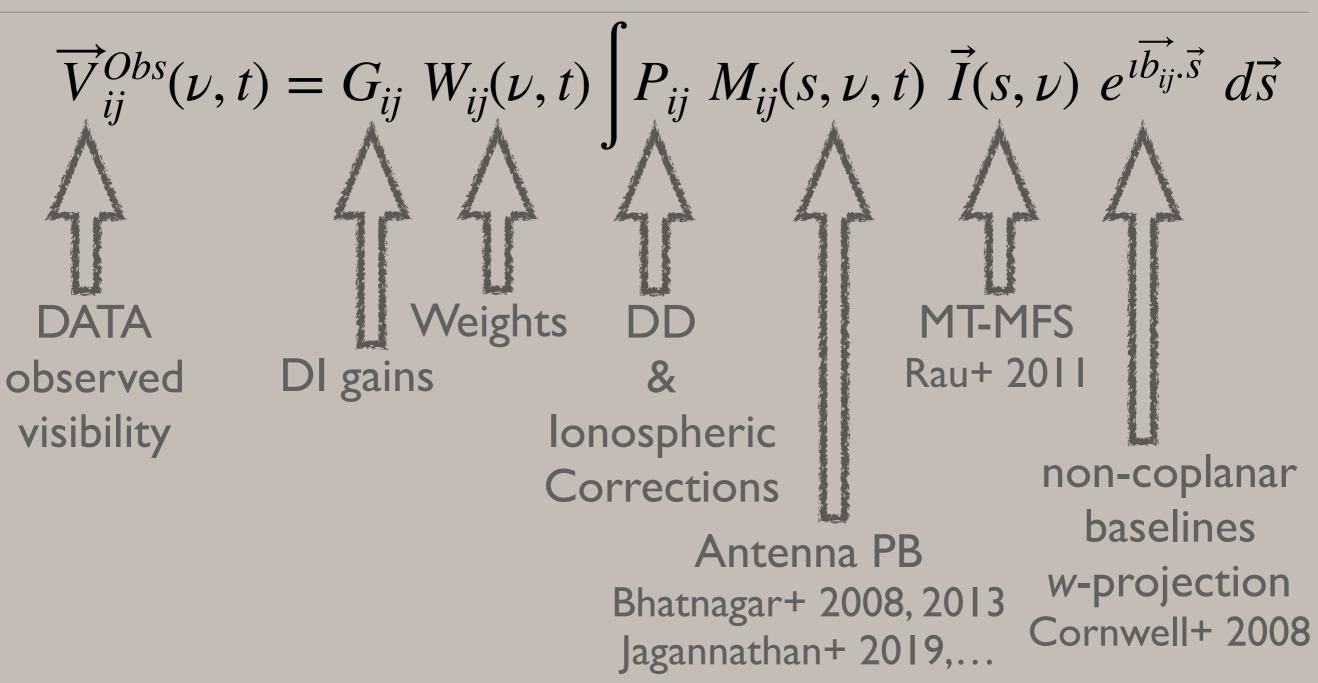
CALIBRATION AND IMAGING

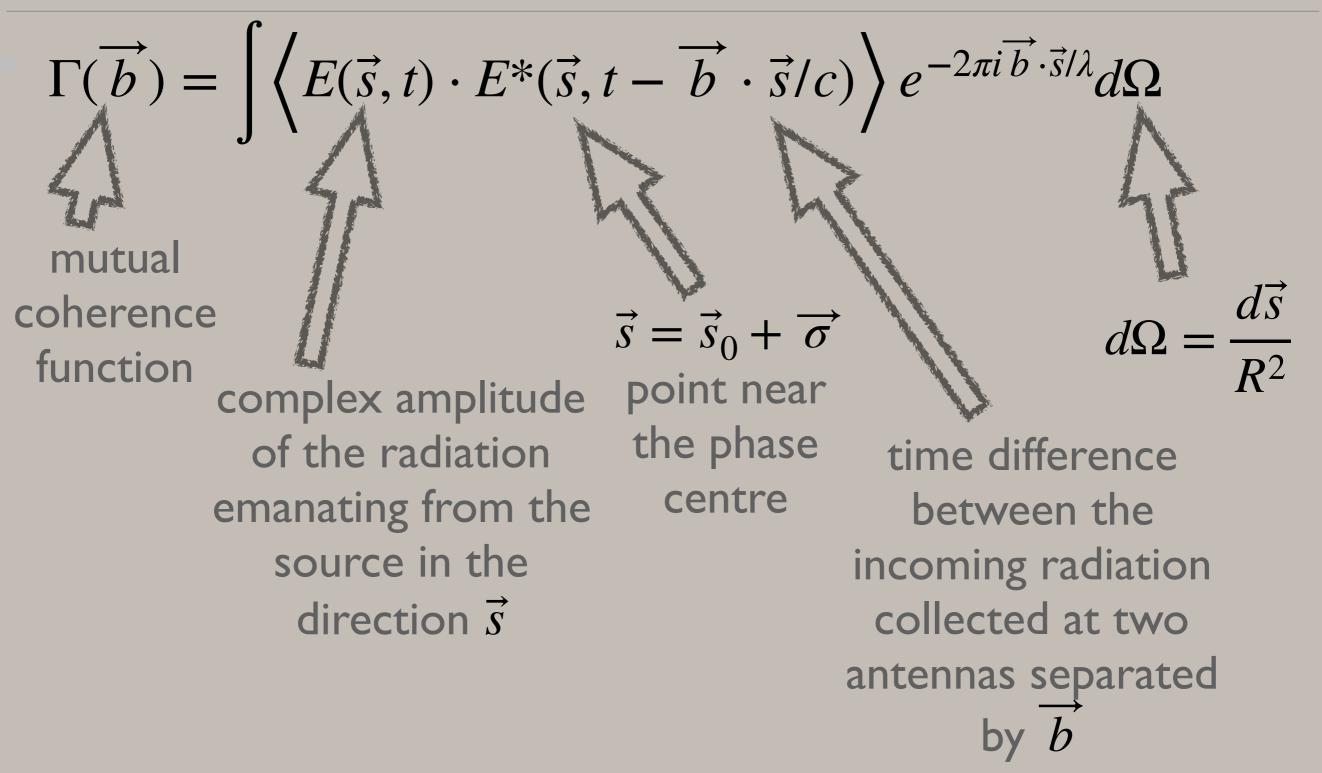
Standard calibration and imaging (DI instrumental effects) w/ DD instrumental + propagation effects correction for *w*-term and for PB image plane correction Fourier plane correction pointing self-calibration Mosaicing w/ advanced image parameterisation multi-scale CLEAN (deconvolution) multi-frequency synthesis (imaging) full polarisation (Stokes) calibration and imaging

 $\overrightarrow{V}_{ij}^{Obs}(\nu,t) = G_{ij} W_{ij}(\nu,t) \left[P_{ij} M_{ij}(s,\nu,t) \overrightarrow{I}(s,\nu) e^{i\overrightarrow{b}_{ij}\cdot\overrightarrow{s}} d\overrightarrow{s} \right]$

$$\overrightarrow{V}_{ij}^{Obs}(\nu,t) = G_{ij} W_{ij}(\nu,t) \int P_{ij} M_{ij}(s,\nu,t) \overrightarrow{I}(s,\nu) e^{i\overrightarrow{b}_{ij}\cdot\overrightarrow{s}} d\overrightarrow{s}$$

$$V(u, v, w) = \int \frac{I(l, m, n)}{n} e^{-2\pi i (ul + vm + w(n-1))} dl dm$$





$$\Gamma(\vec{b}) = \int \left\langle E(\vec{s},t) \cdot E^*(\vec{s},t-\vec{b}\cdot\vec{s}/c) \right\rangle e^{-2\pi i \vec{b}\cdot\vec{s}/\lambda} d\Omega$$

$$V(u,v,w) = \int \frac{I(l,m,n)}{n} e^{-2\pi i (ul+vm+w(n-1))} dl dm$$
(for $w \simeq 0, n \simeq 1$)
$$V(u,v) = \int I(l,m) e^{-2\pi i (ul+vm)} dl dm$$
(this is van-Cittert Zernike theorem)

$$\Gamma(\vec{b}) = \int \left\langle E(\vec{s},t) \cdot E^*(\vec{s},t-\vec{b}\cdot\vec{s}/c) \right\rangle e^{-2\pi i \vec{b}\cdot\vec{s}/\lambda} d\Omega$$
$$V(u,v,w) = \int \frac{I(l,m,n)}{n} e^{-2\pi i (ul+vm+w(n-1))} dl dm$$

Polarised radiation:

 $\overrightarrow{E_i} = [E^r \ E^l]_i^T$

(two nominal orthogonal components of incident electric field are measured at each antenna i)

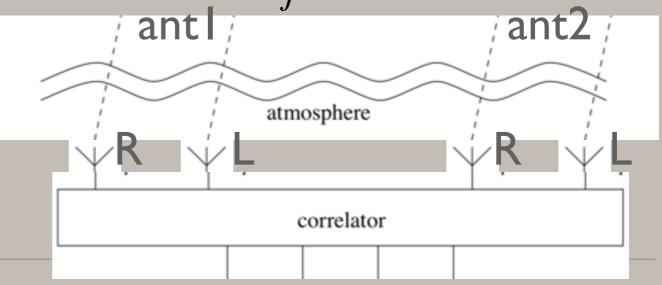
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Polarised radiation:

$$\overrightarrow{E_i} = [E^r \ E^l]_i^T$$

(four cross-correlation products, $\langle \vec{E_i} \otimes \overline{E_i^*} \rangle$ per baseline)

$$\overrightarrow{V_{ij}} = \begin{bmatrix} V^{rr} & V^{rl} & V^{lr} & V^{ll} \end{bmatrix}_{ij}^{T}$$
$$\overrightarrow{I} = \begin{bmatrix} I^{rr} & I^{rl} & I^{lr} & I^{ll} \end{bmatrix}^{T}$$



 $\overrightarrow{E_i} = [E^r \ E^l]_i^T$

(suffers from propagate effects and receiver electronics) (Jones matrices describe this modulation for the electric field incident at each orthogonal pair of feeds $\overline{E_i} = [E^r \ E^l]_i^T$ $DD: J_i^{sky} = [EPF]$ DI: $J_i^{vis} = [GDC]$ (a 2×2 matrix product) (a 2×2 matrix product) complex gains, G, AIPs, E,polar'n leakage, D and PA effects, P and feed config'n, C. tropospheric / ionospheric effects, and Faraday R'n, F.

 $\overrightarrow{E_i} = [E^r \ E^l]_i^T$

(suffers from propagate effects and receiver electronics)

(Jones matrices describe this modulation for the electric field incident at each orthogonal pair of feeds $\vec{E_i} = [E^r \ E^l]_i^T$)

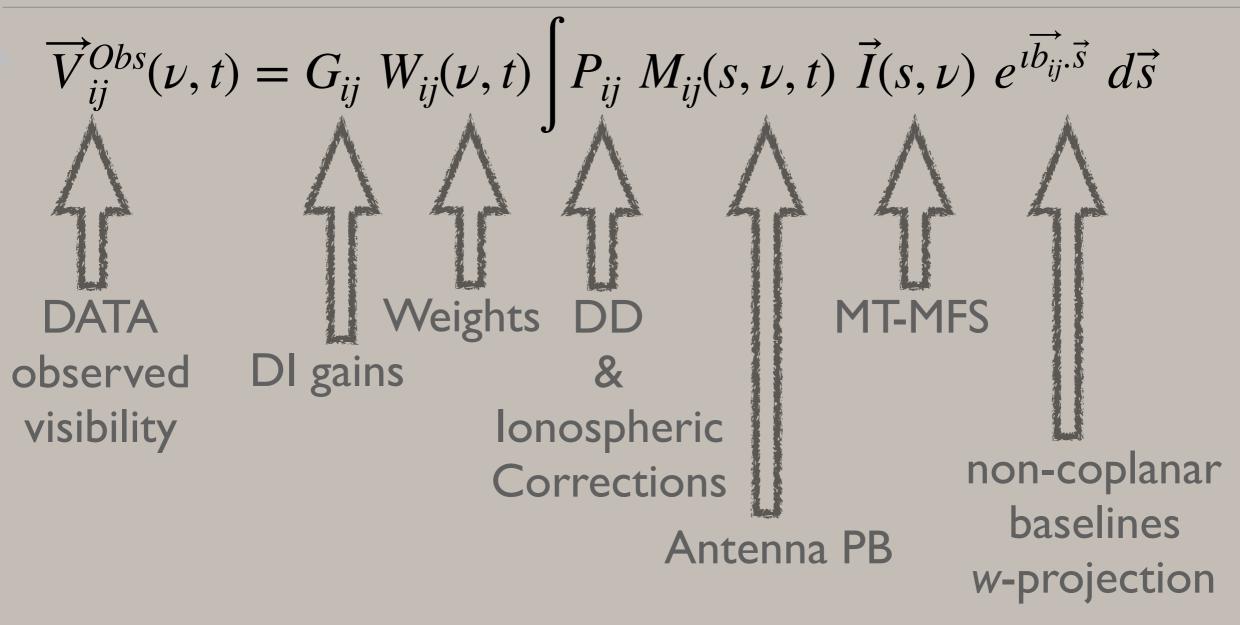
DI:
$$J_i^{vis} = [GDC]$$

DD: $J_i^{sky} = [EPF]$
 $K_{ij}^{\{vis, sky\}} = [J_i \otimes J_j^{\dagger}]^{\{vis, sky\}}$

(effect on each baseline ij is described by the outer-product of these antenna-based Jones matrices, a 4×4 matrix!)

$$\vec{V}_{ij}^{obs} = [K_{ij}^{vis}] \left[[K_{ij}^{sky}] \vec{I}^{sky}(\vec{s}) e^{-2\pi i \vec{b} \cdot \vec{\sigma}/\lambda} d\Omega \right]$$

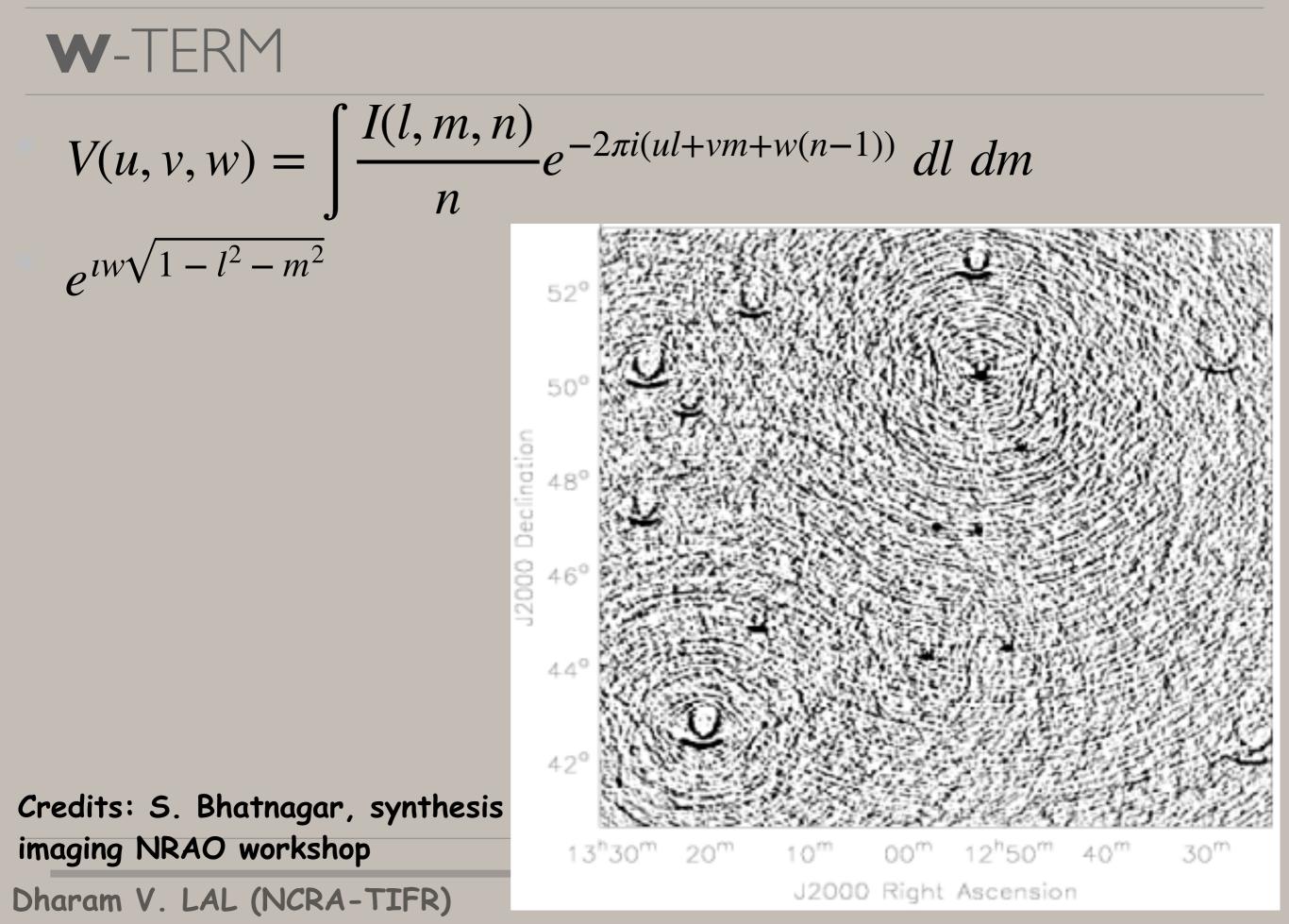




$$\overrightarrow{V}_{ij}^{obs} = [K_{ij}^{vis}] \int [K_{ij}^{sky}] \vec{I}^{sky}(\vec{s}) e^{-2\pi i \vec{b} \cdot \vec{\sigma}/\lambda} d\Omega$$

CALIBRATION AND IMAGING

Standard calibration and imaging (DI instrumental effects) w/DD instrumental + propagation effects correction for *w*-term and for PB image plane correction Fourier plane correction pointing self-calibration Mosaicing w/ advanced image parameterisation multi-scale CLEAN (deconvolution) multi-frequency synthesis (imaging) **full polarisation** (Stokes) calibration and imaging

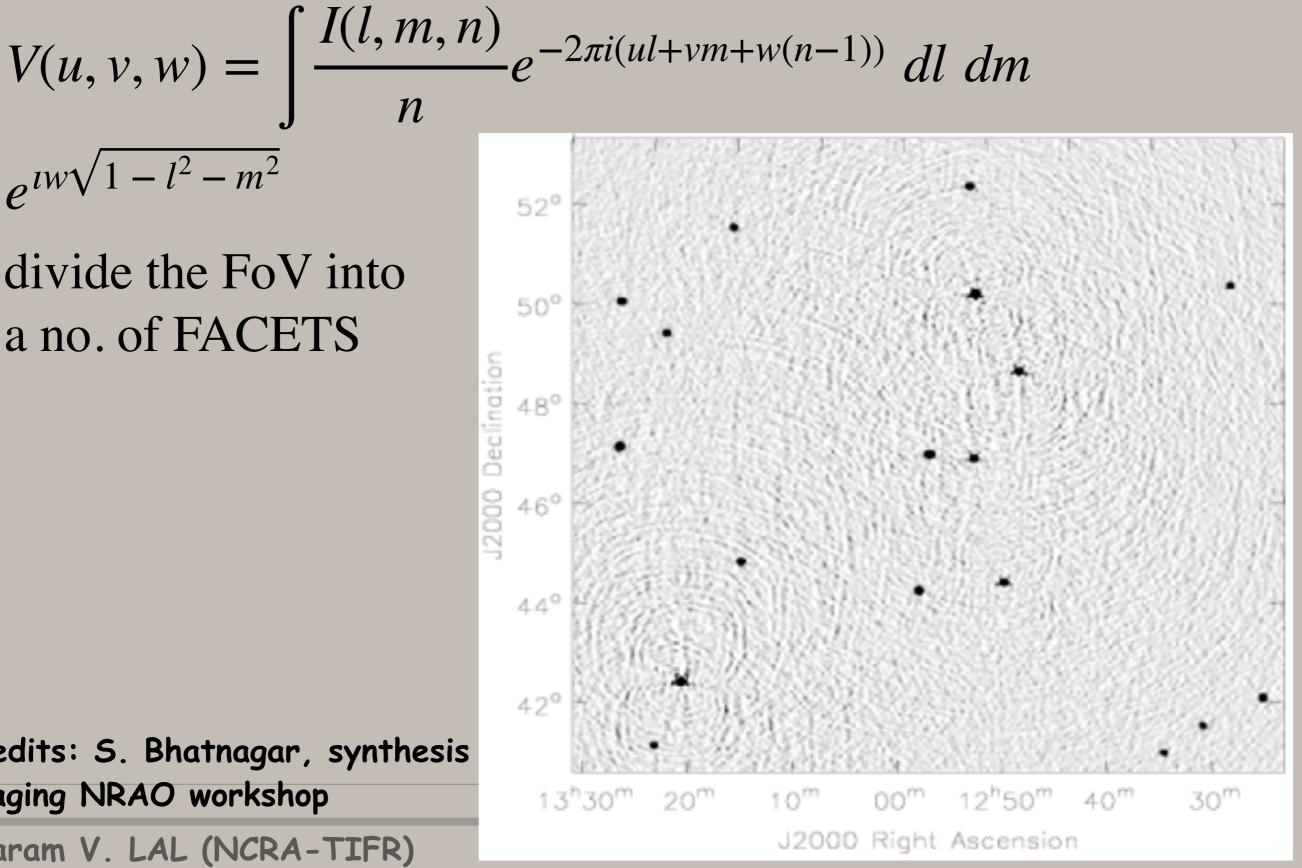


W-TERM

 $e^{lw}\sqrt{1-l^2-m^2}$

divide the FoV into a no. of FACETS

Credits: S. Bhatnagar, synthesis imaging NRAO workshop

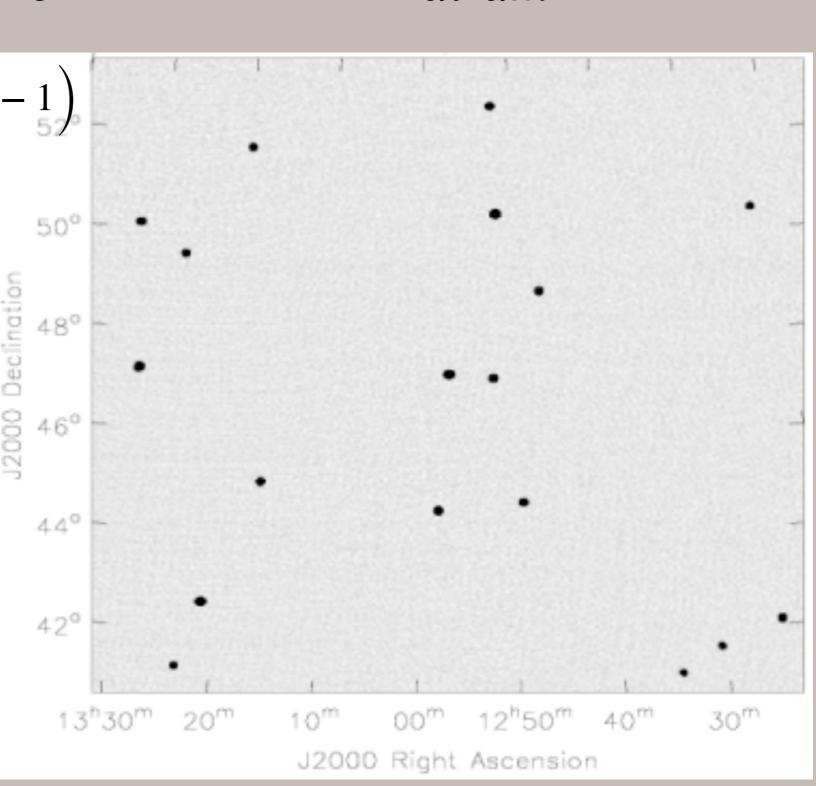


W-TERM

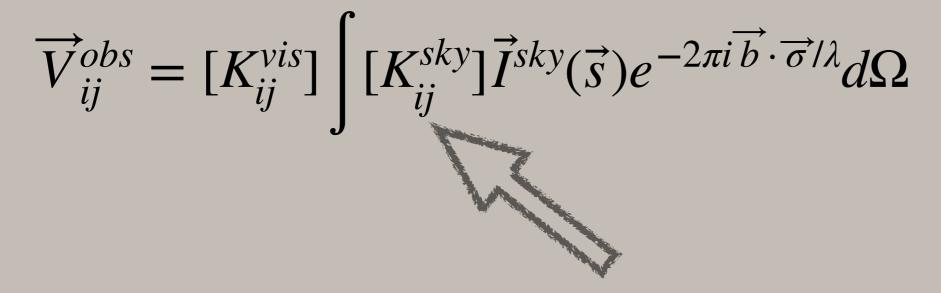
$$V(u, v, w) = \int \frac{I(l, m, n)}{n} e^{-2\pi i (ul + vm + w(n-1))} dl dm$$
$$K_{ii}^{Sky} = e^{w_{ij} \left(\sqrt{1 - l^2 - m^2} - 1\right)}$$

An order-of magnitude faster than FACETing, and

- for the same amount of computing time provides higher DR images.
- Credits: S. Bhatnagar, synthesis imaging NRAO workshop



A-projection



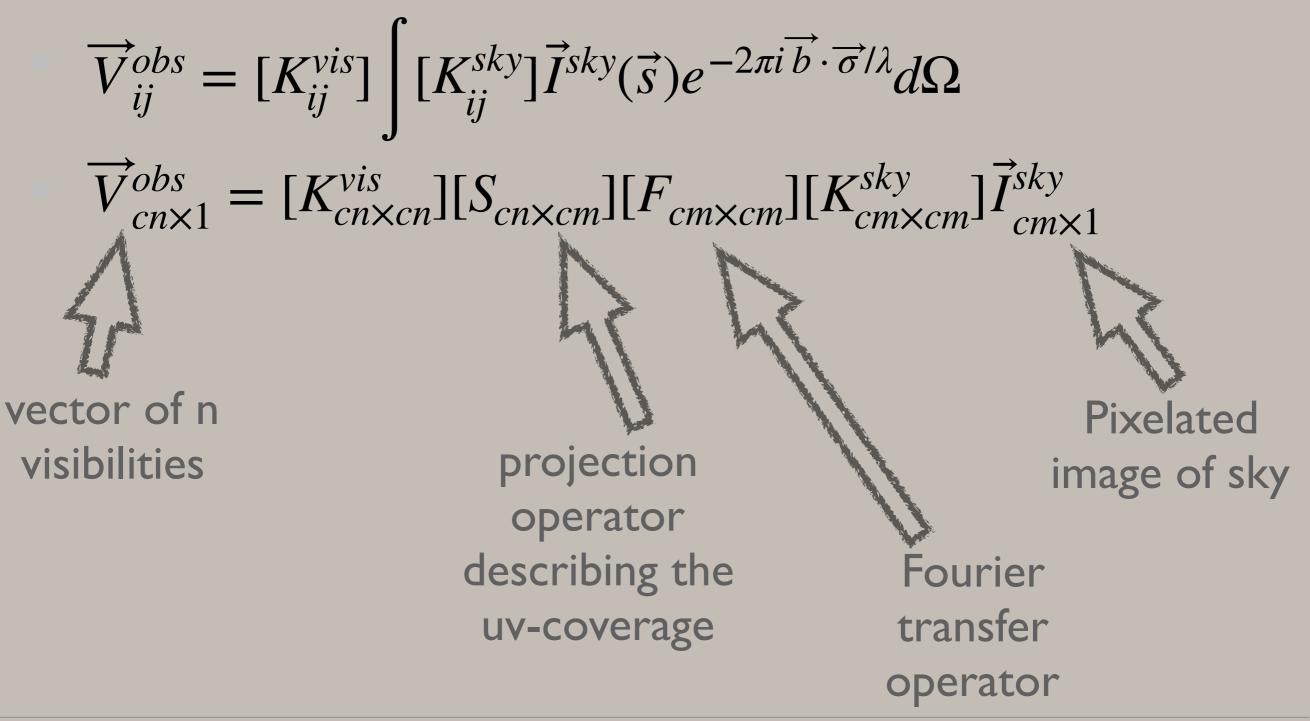
is different for each baseline

Assumption:

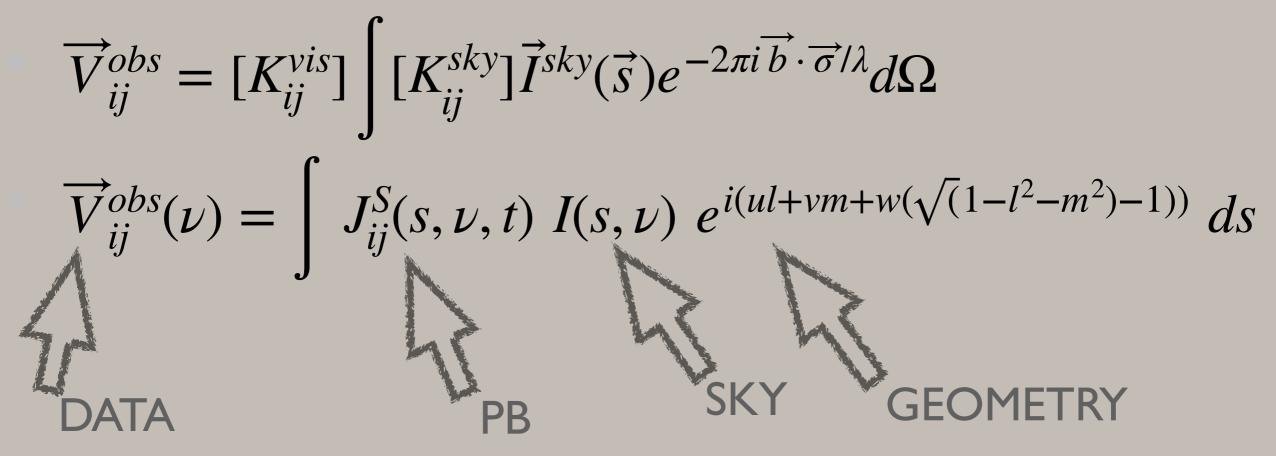
sky is (not) variable, and

Antenna power pattern is (not) changing!

A-projection



A-projection



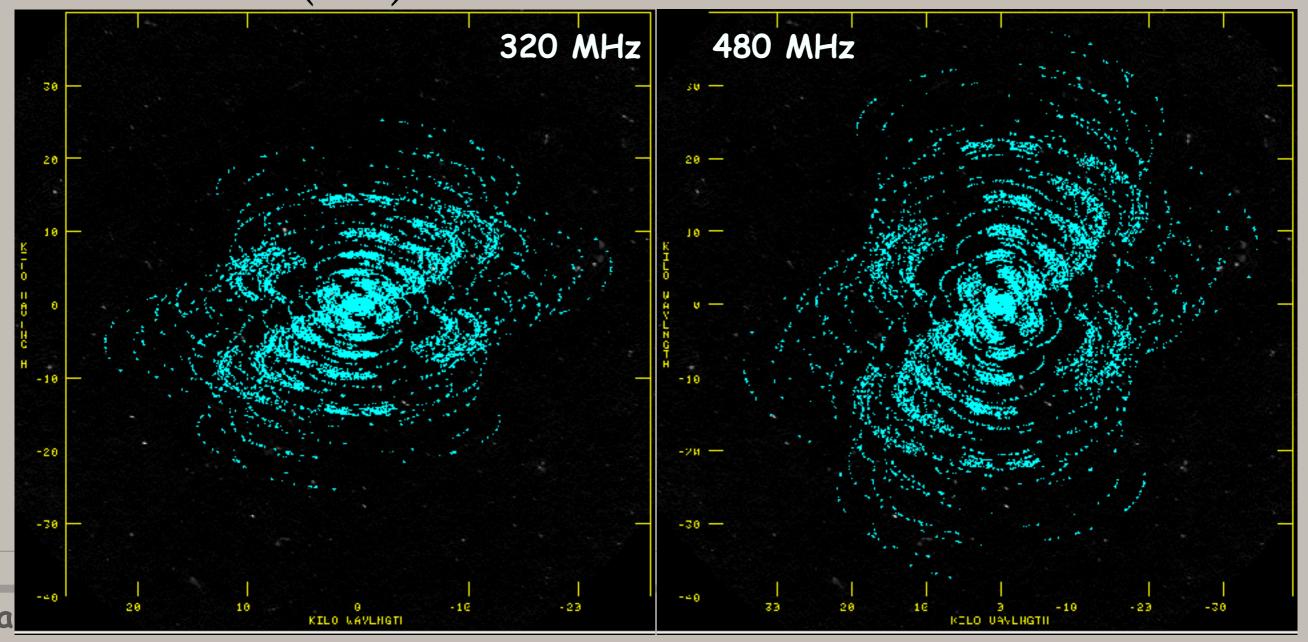
Visibility depends on time and frequency!

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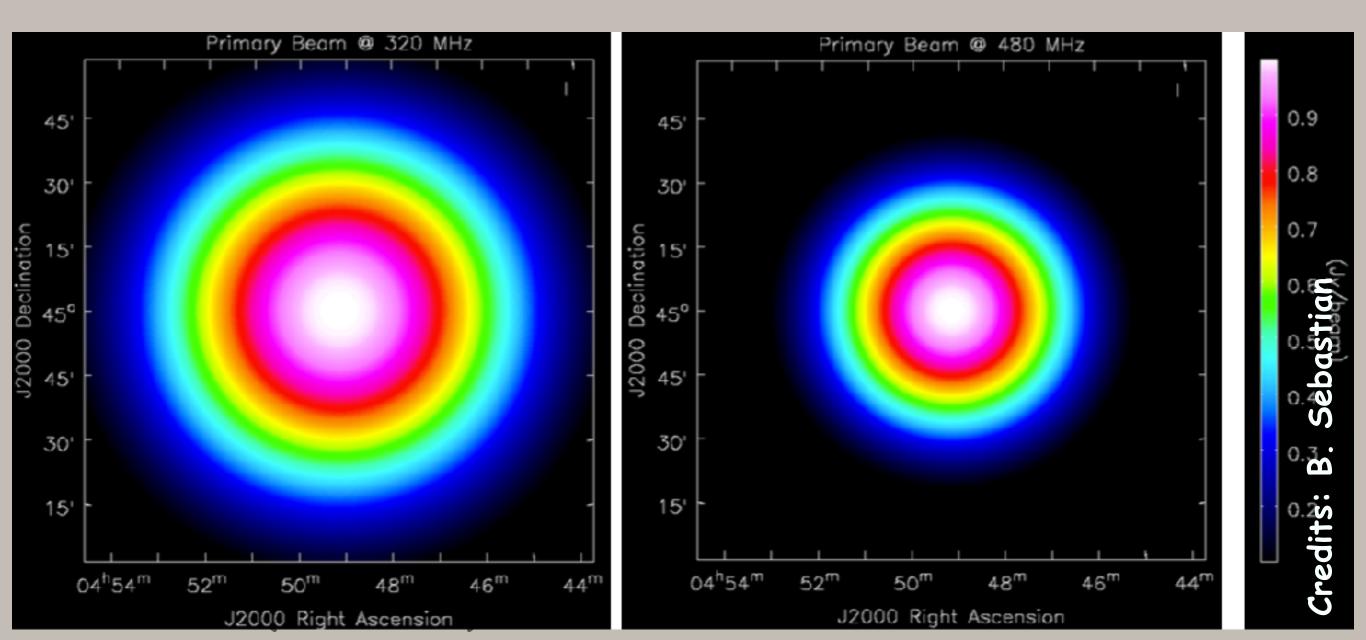
multi-frequency synthesis

$$I_{\nu}^{sky} = I_{\nu0}^{sky} \left(\frac{\nu}{\nu0}\right)^{I_{\alpha}^{sky} + I_{\beta}^{sky} \log(\frac{\nu}{\nu0})}$$



A-projection

 $\overrightarrow{V}_{ij}^{obs}(\nu) = \int J_{ij}^{S}(s,\nu,t) \ I(s,\nu) \ e^{i(ul+\nu m+w(\sqrt{(1-l^2-m^2)-1}))} \ ds$



A-projection

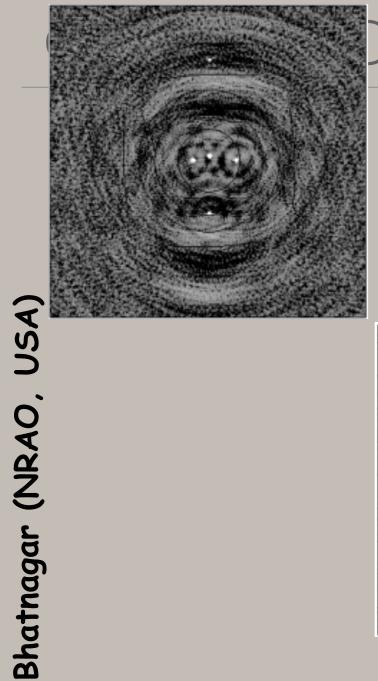
$$\vec{V}_{ij}^{obs}(\nu) = \int J_{ij}^{S}(s,\nu,t) \ I(s,\nu) \ e^{i(ul+\nu m+w(\sqrt{(1-l^2-m^2)}-1))} \ ds$$

multi-frequency synthesis

$$I_{\nu}^{sky} = I_{\nu 0}^{sky} \left(\frac{\nu}{\nu 0}\right)^{I_{\alpha}^{sky} + I_{\beta}^{sky} \log(\frac{\nu}{\nu 0})}$$
$$I_{0} = I_{\nu_{0}}$$
$$I_{1} = I_{\alpha} \times I_{\nu_{0}}$$
$$I_{2} = (I_{\alpha}(I_{\alpha} - 1)/2 + I_{\beta}) \times I_{\nu_{0}}$$

300-400 MHz 26 IA й 39

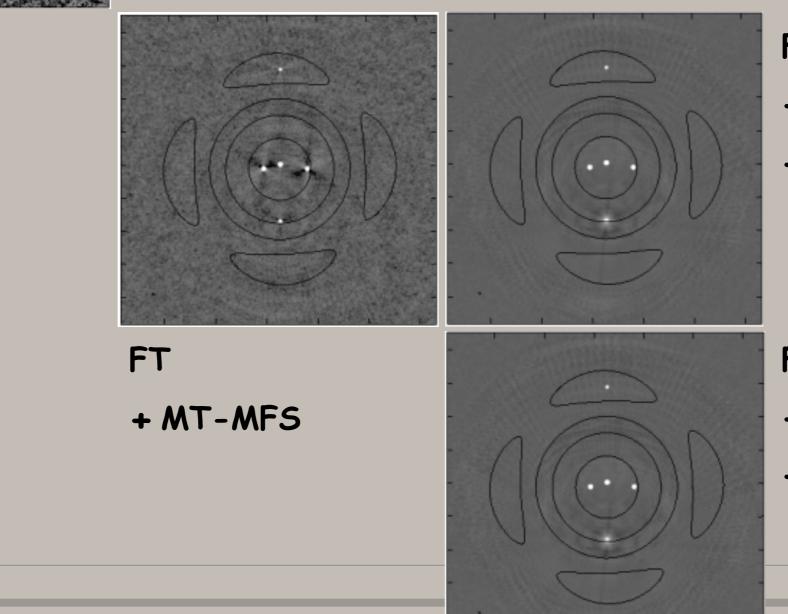
RAS2023: 15 March 2023



DN FOR PB

FT

(standard imaging)



FT

- + MT-MFS
- + A-projection

FT + MT-MFS + WB A-projection

<u>ທ</u>

Credits:

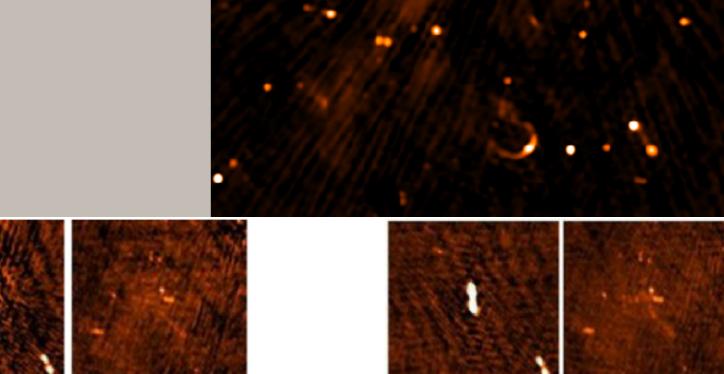
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PEELING: DD CALIBRATION

antenna based gains are determined in the direction of each compact source.

subtract these gains (contribution of compact sources from the observed data using a DFT) and the residual visibilities are imaged again.
drawbacks of peeling...



Dharam V. LAL (NCRA-TIFR)

Credits: H. Intema (Leiden Obs.)

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- Defined by George in 1852
 - Adopted for astronomy by Chandrasehkar in 1947.
 - Not a vector quantity! Deals with power instead of electric field amplitudes.
 - Can be used for partially polarised radiation.
 - The correlator can produce ALL Stokes parameters simultaneously (not so easy in optical astronomy!)

HOW DO WE MEASURE STOKES?

Stokes parameters are the auto-correlation & crosscorrelation products returned from the correlator, but input to the correlator can come from different feed types.

Feeds normally designed to approximate pure linear or circular.

Circular feeds – frequency dependent response adds 90° phase to R for L, so:

I from RR + LL

V from RR - LL

I – total intensity and sum of any two orthogonal polarisations

Q & U – completely specify linear polarisation V – completely specifies circular polarisation $RR = \mathscr{A}(RR)e^{i\psi RR} = I + V$ $LL = \mathscr{A}(LL)e^{i\psi LL} = I - V$ $RL = \mathscr{A}(RL)e^{i\psi RL} = Q + iU$ $LR = \mathscr{A}(LR)e^{i\psi LR} = Q - iU$

I – total intensity and sum of any two orthogonal polarisations

Q & U - completely specify linear polarisation V - completely specifies circular polarisationStokes parameters (as percentages of I) $I = \frac{(RR + LL)}{2} \qquad \qquad \frac{Q}{I} = \frac{\text{Re}(RL + LR)}{RR + LL}$ $\frac{V}{I} = \frac{RR - LL}{RR + LL} \qquad \qquad \frac{U}{I} = \frac{\text{Im}(RL - LR)}{RR + LL}$

I – total intensity and sum of any two orthogonal polarisations

Q & U – completely specify linear polarisation

V – completely specifies circular polarisation

Stokes parameters (as percentages of I)

Is it really that simple?

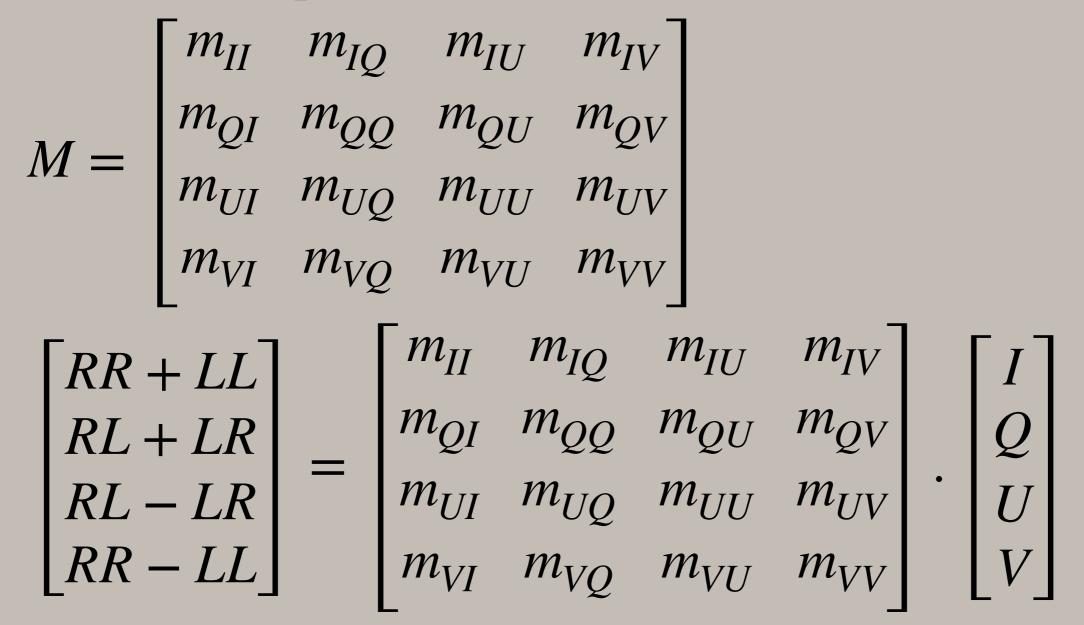
No, there are leakages...

The total intensity can leak into the polarised components (I into {Q,U,V}).

Hans Mueller

MUELLER MATRIX

The leakage of each polarisation into the other can be measured and quantified in a 4×4 matrix (Mueller 1943).



POLARISATION CALIBRATION

- Flux density scale
- $I \Leftrightarrow Q$ leakage
- $I \Leftrightarrow U$ leakage
- $I \Leftrightarrow V$ leakage

Alignment => PA calibration Ellipticity, $Q \Leftrightarrow V$ *RL* phase, $U \Leftrightarrow V$ Constrained using calibrator with known Stokes parameters

POLARISATION CALIBRATION

Flux density scale

- $I \Leftrightarrow Q$ leakage
- $I \Leftrightarrow U$ leakage
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Alignment => PA calibration

Ellipticity, $Q \Leftrightarrow V$

RL phase, $U \Leftrightarrow V$

-

Constrained using calibrator with known Stokes parameters

Need calibrator with known PA

POLARISATION CALIBRATION

Flux density scale

- $I \Leftrightarrow Q$ leakage
- $I \Leftrightarrow U$ leakage
- $I \Leftrightarrow V$ leakage

Alignment => PA calibration

Ellipticity, $Q \Leftrightarrow V$

RL phase, $U \Leftrightarrow V$

Constrained using calibrator with known Stokes parameters

Need calibrator with known PA

> Stokes V ~ 0 for most calibrators so no need to worry too much unless you require very high precision

PUTTING THIS ALL TOGETHER

In the end what we are trying to do is relate products from our correlator to the intrinsic polarised radiation from the source.

So we need to correct the raw correlator outputs for imperfections in the receiver (leakages).

The orientation of the receiver with respect to the telescope structure.

a.k.a. the changing parallactic angle.

Any measured propagation related polarisation effects (e.g. Faraday rotation).

BEAM EFFECTS

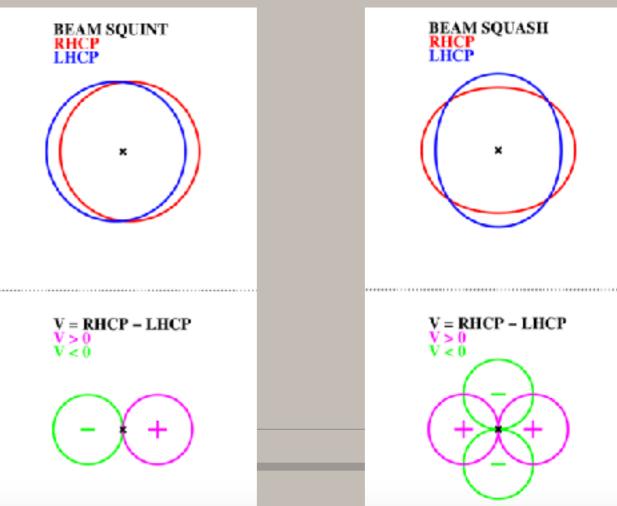
For point sources, all of the previous is fine.

What if the source you are looking at is extended compared to the telescope beam?

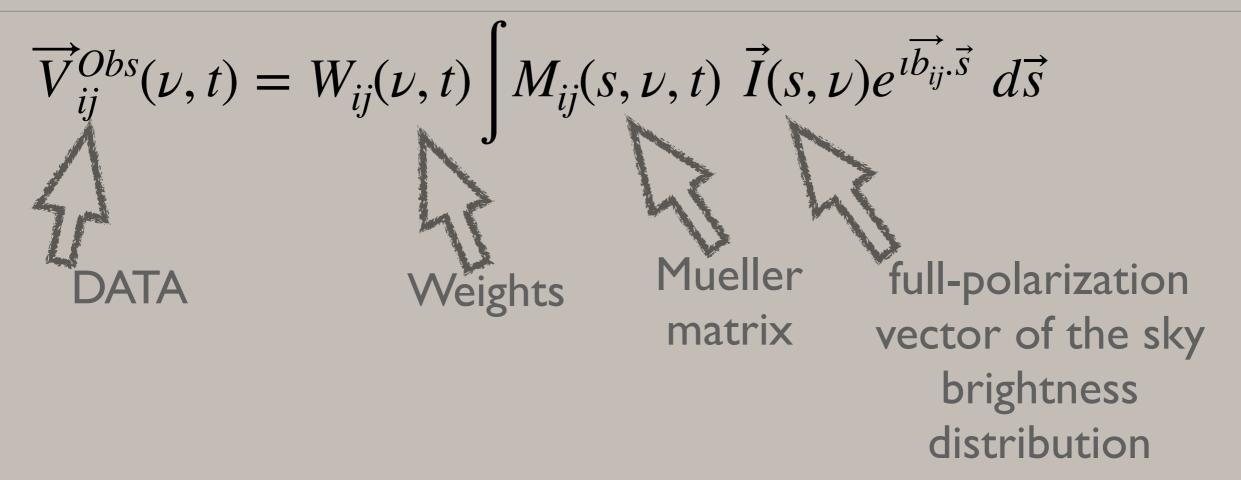
There are instrumental beam effects that can confuse the measurement of extended polarised signals. They are...

Squint

Squash



BEAM EFFECTS



BEAM EFFECTS

$$\overrightarrow{V}_{ij}^{Obs}(\nu, t) = W_{ij}(\nu, t) \int M_{ij}(s, \nu, t) \ \overrightarrow{I}(s, \nu)e^{i\overrightarrow{b}_{ij}\cdot\overrightarrow{s}} \ d\overrightarrow{s}$$

$$M_{ij}(\overrightarrow{s}, \nu, t) = E_i(\overrightarrow{s}, \nu, t) \otimes E_j^*(\overrightarrow{s}, \nu, t)$$

$$\overrightarrow{V}_{ij}^{Obs}(\nu, t) = W_{ij}(\nu, t)\mathscr{F}\left[\left(E_i(\overrightarrow{s}, \nu, t) \otimes E_j^*(\overrightarrow{s}, \nu, t)\right) \cdot \overrightarrow{I}(\overrightarrow{s}, \nu)\right]$$

$$= W_{ij}(\nu, t)\left[A_{ij} \star \overrightarrow{V}_{ij}\right]$$
where, $A_{ij} = A_i \otimes A_j^*$

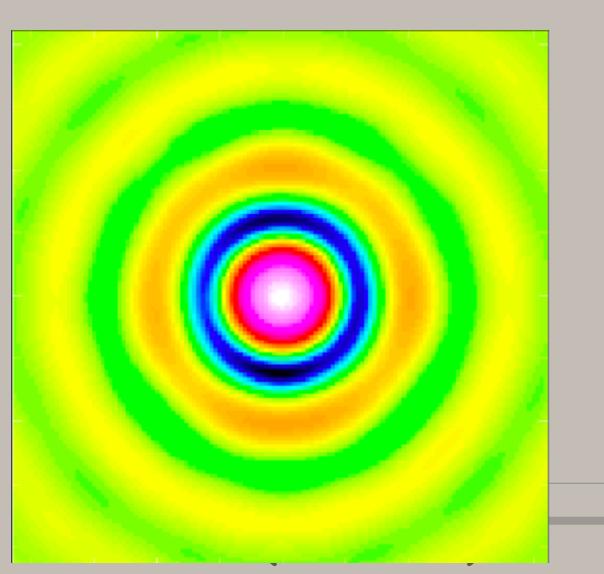
$$AIPs \text{ for two antenna}$$

$$correlator$$

Dharam V. LAL (NCRA-TIFR)

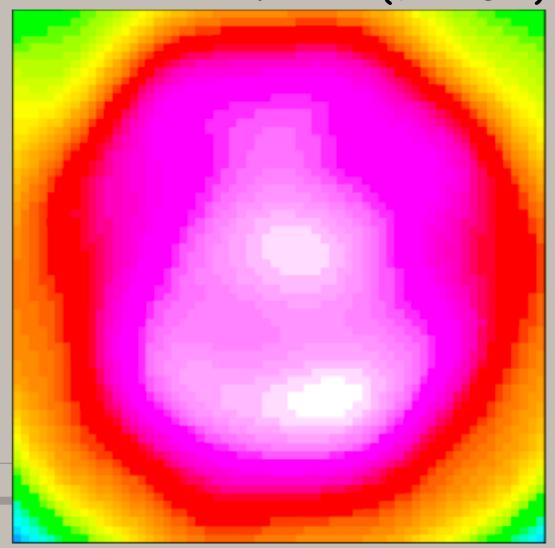
APERTURE ILLUMINATION

Holography data: MeerKAT Obtained by Fourier Transforming the PB Holography measurements





Credits: S. Sekhar (UCT-IDIA)



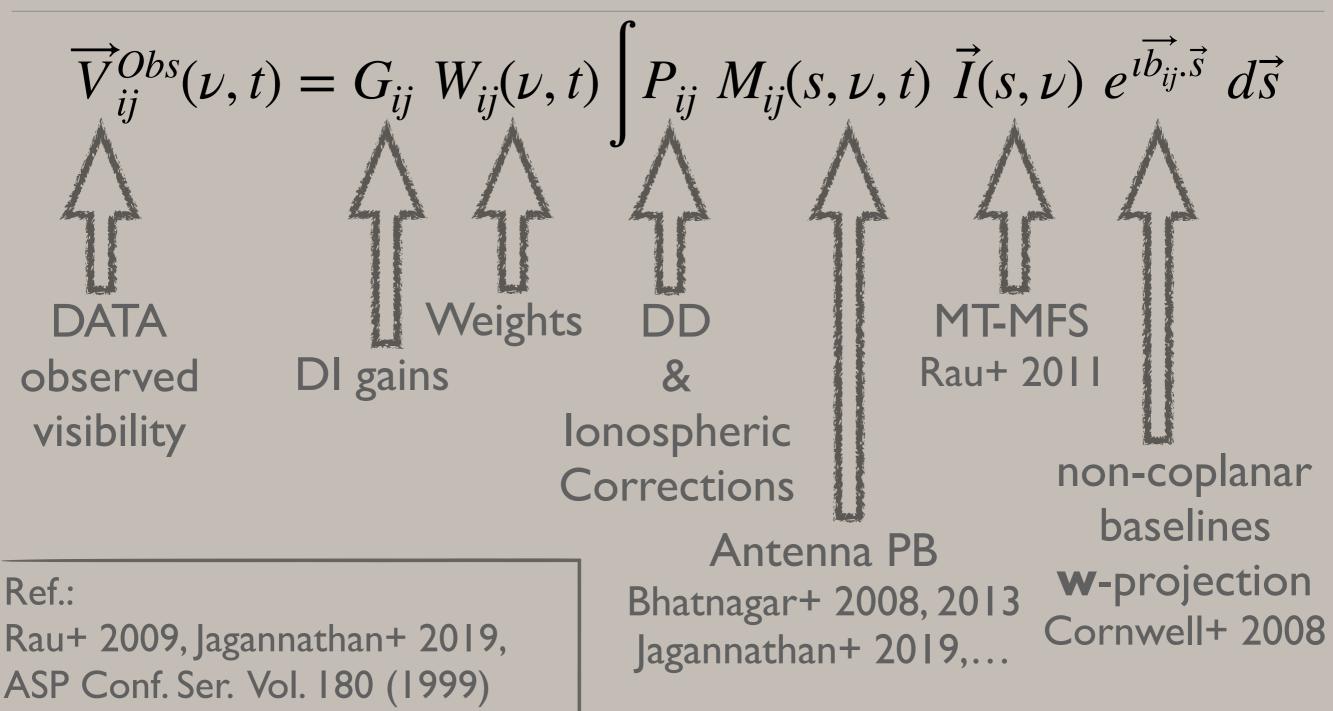
A-TO-Z SOLVER

Use Zernike polynomials to directly model the complex aperture

it is a natural domain to model optical aberrations that cause PB weirdness

(Telescope agnostic - does not require ray traced model for different antennas/telescopes, only Holography)

Aperture size is fixed, independent of number of measured sidelobes.



Thank you all for your attention!